## Theoretical Physicists' Biases Meet Machine Learning

Using and finding functional bias in ML for mathematical physics systems
29.10.2021, IAIFI Colloquium


Our purpose in theoretical physics is not to describe the world as we find it, but to explain - in terms of a few fundamental principles - why the world is the way it is.

Steven Weinberg


## Can ML achieve this? [requiring explainable Al]

## If yes, which NEW physics can we reveal?

## Content

Theoretical physics problems made for ML: understanding high-dimensional data
I. Efficient solutions to PDEs (in mathematical physics) with ML

Key: Using domain knowledge/bias in ML ansatz
Example: Numerical Calabi-Yau metrics
II.

How to extract domain knowledge/biases with ML (e.g. what are the symmetries of a system)

## Why ML and physics? <br> ML can overcome curses of dimensionality

- Efficient functional biases can overcome this curse of dimensionality, e.g. utilising symmetries of your data

Translation invariance: CNNs



- Such functional biases (e.g. symmetries) are at the heart of all physics models

Moduli dependent Calabi-Yau and SU(3)-structure metrics from Machine Learning based on (2012.04656), in collaboration with:


Lara Anderson


Mathis Gerdes


James Gray


Nikhil Raghuram


Fabian Ruehle

Finding symmetries and integrable structures of physical systems
and based on (2104.14444, 2103.07475, 2003.13679, 2002.05169), in collaboration with:


Marc Syvaeri


Dieter Lüst

How to improve our knowledge of EFTs in string theory with the metric (non-holomorphic quantities)?

## Metrics with ML

## Metrics matter

- The metric is key in any extra-dimensional physics model

$$
S=\int_{M_{4+D}} d^{4+D} x \sqrt{-\operatorname{det} g_{4+D}} R\left(\underset{4+D}{g_{4+D}}\right)
$$



- String compactifications are no exception to this. For instance:

1. Matter kinetic terms (soft-terms, cf. 0906.3297)
2. Moduli potential (D3-brane inflation [probing directly CY-moduli space])
3. Massive string spectrum

## Signatures of Quantum Gravity

## Metrics in the EFT

How to distinguish these signatures from some bottom-up BSM model?

Characteristic features in the EFTs of theories with extra dimensions?

$$
\mathscr{L}_{\text {moduli }}=k(\phi)(\partial \phi)^{2}+V(\phi)
$$

Understand the string theory EFT better

Is this picture true?

Stringy $k(\phi)$-space

$$
S=\int_{M_{4+D}} d^{4+D} x \sqrt{-\operatorname{det} g_{4+D}} R\left(g_{4+D}\right)
$$



## Which Metrics?

## 6D metrics relevant for string theory

- String Theory EOM for 4D $\mathcal{N}=1$ Minkowski vacua require a Ricci-flat Kähler metric (Candelas, Horowitz, Strominger, Witten 1985)
- Which compact spaces do exist with a Ricci-flat Kähler metric?


## Calabi-Yau manifolds

(Example today: Quintic hypersurface in $\mathbb{P}^{4}$ )

- Yau (1977) showed the existence of such a unique Ricci-flat Kähler metric, but without explicit constructions.
- One definition of CY-threefold: complex threefold admitting a nowhere vanishing real two-form J , and a complex three form $\Omega$ such that:

$$
J \wedge \Omega=0, J \wedge J \wedge J=\frac{3 i}{4} \Omega \wedge \bar{\Omega}, d J=0, d \Omega=0
$$

## Which Metrics?

## 6D metrics relevant for string theory

- The metric is given as $i g_{a \bar{b}}=J_{a \bar{b}}$
- Simplest examples: complete intersection manifolds in projective spaces

Quintic hypersurface in $\mathbb{P}^{4}$ :
$p_{\psi}(\vec{z})=\sum_{i=0}^{d+1} z_{i}^{d+2}+\psi \prod_{i=0}^{d+1} z_{i}=0$

Algebraic metrics:
$K=1 / 2 \pi \ln (\mathbf{k})$
$\mathbf{k}=\sum_{\alpha, \bar{\beta}=0}^{N_{k}} s_{\alpha}(\vec{z}) H_{\alpha \bar{\beta}} \bar{s}_{\bar{\beta}}(\overrightarrow{\bar{z}})$
$g_{a \bar{b}}=\partial_{a} \bar{\partial}_{\bar{b}} K=\frac{1}{2 \pi} \frac{\mathbf{k} \mathbf{k}_{a \bar{b}}-\mathbf{k}_{a} \mathbf{k}_{\bar{b}}}{\mathbf{k}^{2}}$

## Which Metrics?

## Functional bias: algebraic metrics

- Idea: Generalised Fubini Study metrics can approximate the metric of our choice

$$
\begin{aligned}
& K=1 / 2 \pi \ln (\mathbf{k}), \mathbf{k}=\sum_{\alpha, \bar{\beta}=0}^{N_{k}} s_{\alpha}(\vec{z}) H_{\alpha \bar{\beta}} \bar{s}_{\bar{\beta}}(\overrightarrow{\bar{z}}) \\
& g_{a \bar{b}}=\partial_{a} \bar{\partial}_{\bar{b}} K=\frac{1}{2 \pi} \frac{\mathbf{k k}_{a \bar{b}}-\mathbf{k}_{a} \mathbf{k}_{\bar{b}}}{\mathbf{k}^{2}}
\end{aligned}
$$

- Embedding into larger projective space (Kodaira embedding): $s_{\alpha}(\vec{z})$ polynomials in $z_{a}$.
- These metrics provide "basis" of Kähler metrics on X. (Tian: such Kähler potentials are dense in the space of Kähler potentials)

Quintic hypersurface in $\mathbb{P}^{4}$ :
$p_{\psi}(\vec{z})=\sum_{i=0}^{d+1} z_{i}^{d+2}+\psi \prod_{i=0}^{d+1} z_{i}=0$

Algebraic metrics:
$K=1 / 2 \pi \ln (\mathbf{k})$
$\mathbf{k}=\sum_{\alpha, \bar{\beta}=0}^{N_{k}} s_{\alpha}(\vec{z}) H_{\alpha \bar{\beta}} \bar{s}_{\bar{\beta}}(\overrightarrow{\bar{z}})$
$g_{a \bar{b}}=\partial_{a} \bar{\partial}_{\bar{b}} K=\frac{1}{2 \pi} \frac{\mathbf{k k}_{a \bar{b}}-\mathbf{k}_{a} \mathbf{k}_{\bar{b}}}{\mathbf{k}^{2}}$

## Metrics are hard without ML

## 6D metrics relevant for string theory

- Finite distance methods "fail" (Headrick, Wiseman 2009)
- Spectral methods simplify, but they are currently inefficient:

1. Single point in moduli space
2. High accuracies become expensive
(Donaldson, Braun, Belidze, Douglas, Ovrut, Karp, Cui, Gray, Lukic, Ashmore, He; Kachru, Tripathy, Zimet; Headrick and Nasar)

- How about non-Kähler solutions?
- Target on a practical level: metric with reasonable accuracy for one string compactification $\sim \mathrm{O}$ (1 day) [impossible with non ML algorithms]

$\mathrm{k} \sim$ accuracy of spectral resolution

Time to check accuracy of solution $\sigma$


## Can Machine Learning help?

$$
J \wedge J \wedge J \sim \operatorname{det} g
$$

## Which metric?

$$
\Omega=\frac{1}{\partial p_{\psi}(\vec{z}) / \partial z_{b}} \bigwedge_{c=1, \ldots, d} d z
$$

What is the optimisation problem
Monge-Ampere Loss (different metrics)

1. Ricci-flatness: (Induced FS is not Ricci-flat):

Ricci tensor: $R_{i \bar{\jmath}}=-\partial_{i} \bar{\partial}_{\bar{\jmath}} \log \operatorname{det} g$
Cheaper alternative (less derivatives) via Monge-Ampere equation:

$$
J \wedge J \wedge J=\kappa \Omega \wedge \bar{\Omega} \quad \rightarrow \quad \mathscr{L}_{\mathrm{MA}}=\frac{1}{\int_{X} \Omega \wedge \bar{\Omega}} \int_{X}\left|1-\frac{1}{\kappa} \frac{J^{3}}{\Omega \wedge \bar{\Omega}}\right|
$$

2. Kählerity:

$$
\begin{aligned}
& d J=0 \quad \leftrightarrow \quad g_{i \bar{\jmath}, k} d z_{i} \wedge d \bar{z}_{\bar{\jmath}} \wedge d z_{k}=0=g_{i \bar{\jmath}, \bar{k}} d z_{i} \wedge d \bar{z}_{\bar{\jmath}} \wedge d \bar{z}_{\bar{k}} \\
& c_{i j k}=g_{i \bar{\jmath}, k}-g_{k \bar{\jmath}, i}=0 \quad \rightarrow \quad \mathscr{L}_{\mathrm{d} J}=\sum_{i, j, k}| | \operatorname{Re}\left(c_{i j k}\right)| |_{n}+\left|\left|\operatorname{Im}\left(c_{i j k}\right)\right|\right|_{n}
\end{aligned}
$$


3. Well defined across different coordinate patches:

$$
g^{(j)}=T_{i j} \cdot g^{(i)} \cdot T_{i j}^{\dagger}, \quad T_{i j}=\partial \vec{z}^{(i)} / \partial \vec{z}^{(j)} \quad \rightarrow \quad \mathscr{L}_{\text {Transition }}=\frac{1}{d} \sum_{k, j}| | g_{\mathrm{NN}}^{(k)}(\vec{z})-T_{j k}(\vec{z}) \cdot g_{\mathrm{NN}}^{(j)}(\vec{z}) \cdot T_{j k}^{\dagger}(\overrightarrow{\vec{z}})| |_{n}
$$

## Our experiments

## Overview on what we get to work

- Supervised learning of Kähler potential (data from running spectral algorithms) Improvement: moduli dependence of metric
- Unsupervised learning of Kähler potential (using energy functionals measuring deviation from Ricci-flatness) Improvement: moduli dependence of metric and efficiency (no running of spectral methods)

(no running of spectral methods)
- Unsupervised learning of metric directly (perturbation of Fubini study metric)
- Metric networks to go beyond Calabi-Yau: here SU(3) structure manifolds, i.e. more general string backgrounds


## Learning H

## Optimising with $\sigma$ (no Donaldson)



- $\mathrm{k}=6$ (42025 components in H), sampling fast and always using new points

$$
\sigma=\frac{1}{\int_{X} \Omega \wedge \bar{\Omega}} \int_{X}\left|1-\frac{1}{\kappa} \frac{J^{3}}{\Omega \wedge \bar{\Omega}}\right|
$$



Algebraic metrics:

$$
\begin{aligned}
& K=1 / 2 \pi \ln (\mathbf{k}) \\
& \mathbf{k}=\sum_{\alpha, \bar{\beta}=0}^{N_{k}} s_{\alpha}(\vec{z}) H_{\alpha \bar{\beta}} \bar{s}_{\bar{\beta}}(\overrightarrow{\bar{z}}) \\
& g_{a \bar{b}}=\partial_{a} \bar{\partial}_{\bar{b}} K=\frac{1}{2 \pi} \frac{\mathbf{k} \mathbf{k}_{a \bar{b}}-\mathbf{k}_{a} \mathbf{k}_{\overline{\bar{b}}}}{\mathbf{k}^{2}}
\end{aligned}
$$

Quintic hypersurface in $\mathbb{P}^{4}$ :

$$
p_{\psi}(\vec{z})=\sum_{i=0}^{d+1} z_{i}^{d+2}+\psi \prod_{i=0}^{d+1} z_{i}=0
$$

## Beyond Calabi-Yau Metrics with ML

- Approach of learning metric directly allows to search for metrics with different properties
- Philosophy: modified loss functions, additionally learned outputs.
- Augment the landscape of metrics to G2 and SU(3) structure manifolds? Phenomenologically necessary, otherwise missing large parts of string theory constructions; unexplored mathematical structures.
- Example SU(3) structure manifolds (simple example works)



## Neural networks for differential equations <br> Going beyond CY metrics

- Can NN give efficient approximations to PDE solutions?
- Motivation beyond universal approximation scheme (NN can be shown to give good and accurate predictions to PDEs):
- Solutions to high-dimensional Schrödinger equations (Rupp, Tkatchenko, Müller, von Lilienfeld 2012, ...)
- Black-Scholes PDE (Grohs, Hornung, Jentzen, von Wurstemberger 2018, ...)
- Approximation rates of NNs to solutions of PDEs (Kutyniok, Petersen, Raslan, Schneider 2019, ...)
- SimDL workshop at ICLR 2021

What to do when we do not have domain knowledge?
Can we use Al to identify the correct domain knowledge?

Underlying questions:

## Are we missing mathematical/physical structures?

Can we find such structures with ML and then use them?

## In Chemistry pre 1869?

## Learning atoms for materials discovery

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## Significance

Motivated by the recent achievements of artificial intelligence (Al) in linguistics, we design Al to learn properties of atoms from materials data on its own. Our work realizes
knowledge representation of atoms via computers and could serve as a foundational step toward materials discovery and design fully based on machine learning.

## In Particle Physics pre ~ 60s/70s?



$$
\begin{aligned}
& \mathcal{L}=-\frac{1}{4} F_{N \nu} F^{N N} \\
& +i \bar{X} \nmid+h_{c} \\
& +x_{i} y_{i j} x_{j} \phi_{+1} \\
& +\left|D_{\mu} \phi\right|^{2}-V(\phi)
\end{aligned}
$$

## Which tools do we need to make such discoveries with ML in the 2020s?



Pattern in Calabi-Yau data



CY-metrics


Finding mathematical structures to describe systems more efficiently

Our approach: Symmetries, Dualities, and Integrability

Why care for ML systems? Symmetries, dualities and integrability are standard structures used in physical systems which make your life easier (parameter inference, predictions from functional bias)
$\rightarrow$ good functional bias

## Symmetries from embedding layer

## How to search for symmetries?

The problem


1. How to find invariances?

$$
f(\phi)=f(\tilde{\phi})
$$

2. Which symmetry is behind such an invariance?


## How to search for symmetries? <br> Embedding in deep layer



We need: group input with the same meaning together
Word2Vec does it:
(England - London = Paris - France)
[1301.3781, used for re-discovering periodic table 1807.05617,
 classifying scents of molecules 1910.10685]


## How to determine the symmetry?

Connected points in input space:

Which symmetry?


## Other Examples?

Determine generator connecting points in (sub)-space:

$$
p^{\prime}=p+\epsilon_{a} T^{a} p
$$

Repeat multiple times (covering all sub-spaces) and perform PCA on generators:




# Symmetries from data (samples of phase space) 

## Simulations and physics bias

- The correct functional expressivity is key (vision: CNNs; geometric deep learning). Example for prediction of trajectories:




## AI and Physics for Simulations

Physics Bias helps for predictions!


## Can we learn more structures from samples of phase space?

## More structures from neural networks?

- If we can train NNs to find the Hamiltonian of a system, can we use it to learn other interesting structures?
- Symmetries of the system? E.g. via canonical transformations (cyclic coordinates reveal conserved quantities)
- How does this work? 2 key steps:

1. Formulate your physics search problem as an optimisation problem.
2. Make sure it's learnable for your architecture.

- Good news for analytic understanding of numerical approximations: most physics functions are simple (AI Feynman [Udrescu, Tegmark 1905.11481])
- Interesting side effect: quantify how much these structures help in predicting dynamics


## Al for Simulations - Symmetries

## Introducing physicists' bias

SCNNs: We cannot only learn the Hamiltonian but also the symmetries by enforcing canonical coordinates


Modified Losses:

$$
0=\dot{F}_{k}(p, q)=\left\{H(p, q), F_{k}(p, q)\right\}
$$

Additional constraint on motion (not just energy conservation), i.e. motion takes place on hyper-surface in phase space


## Al for Simulations - Symmetries <br> Introducing physicists' bias

SCNNs: We cannot only learn the Hamiltonian but also the symmetries by enforcing canonical coordinates


Modified Losses for canonical coordinates:

- Hamilton equations:

$$
\begin{aligned}
& \dot{P}_{i}(p, q)=-\frac{\partial H(p, q)}{\partial Q_{i}(p, q)}=0 \quad \text { and } \quad \dot{Q}_{i}(p, q)=\frac{\partial H(p, q)}{\partial P_{i}(p, q)} \\
& \left\{P_{i}, Q_{j}\right\}=\delta_{i j} \quad \text { and } \quad\left\{P_{i}, P_{j}\right\}=\left\{Q_{i}, Q_{j}\right\}=0
\end{aligned}
$$

- Poisson algebra:



## Benefits from Physicists' Bias

- Conserved quantities interpretable:

$$
\begin{aligned}
& P_{c_{1}}=-4.2 p_{x_{1}}-4.2 p_{x_{2}}-1.3 p_{y_{1}}-1.3 p_{y_{2}}, P_{c_{2}}=-0.9 p_{x_{1}}-0.9 p_{x_{2}}-3.2 p_{y_{1}}-3.2 p_{y_{2}} \\
& L=-1.1 q_{x_{1}} p_{y_{1}}+0.9 q_{x_{1}} p_{y_{2}}+0.9 q_{x_{2}} p_{y_{1}}-1.0 q_{x_{2}} p_{y_{2}}+1.0 q_{y_{1}} p_{x_{1}}-0.9 q_{y_{1}} p_{x_{2}}-0.9 q_{y_{2}} p_{x_{1}}+1.0 q_{y_{2}} p_{x_{2}}
\end{aligned}
$$

- Using learned conserved quantities helps in predicting trajectories



## Can we search for new mathematical/physical structures?

## Symmetries $\rightarrow$ Integrability

## Integrability

## A lightning overview

- Additional constraint $F_{k}$ on motion:

$$
0=\dot{F}_{k}=\left\{H, F_{k}\right\}
$$

How many $F_{k}$ can there be?

- System (2n dimensional) integrable iff: n independent, everywhere differentiable integrals of motion $F_{k}$ (in involution).
- Alternatively search for Lax pair:

$$
\dot{L}=[L, M]
$$

s.t. eom are satisfied. Conserved quantities via:

$$
F_{k}=\operatorname{tr}\left(L^{k}\right)
$$

(additional condition for $\left\{F_{k}, F_{j}\right\}=0$ )

## Example: Harmonic Oscillator

- Hamiltonian and EOM:

$$
H=\frac{1}{2} p^{2}+\frac{\omega^{2}}{2} q^{2} ; \quad \dot{q}=p, \dot{p}=-\omega^{2} q
$$

- Lax pair:

$$
L=a\left(\begin{array}{cc}
p & b \omega q \\
\frac{\omega}{b} q & -p
\end{array}\right), \quad M=\left(\begin{array}{cc}
0 & \frac{b}{2} \omega \\
-\frac{\omega}{2 b} & 0
\end{array}\right)
$$

- Conserved quantities:

$$
\begin{aligned}
& F_{1}=2 \lambda \\
& F_{2}=2 \lambda^{2}+4 H \\
& F_{3}=2 \lambda^{3}+12 \lambda H \quad \lambda \ldots \text { spectral parameter }
\end{aligned}
$$

## Integrability

## We need some deus ex machina moment

Having a Lax pair formulation of integrability is very convenient, but

- inspiration is needed to find it,
- its structure is hardly transparent,
- it is not at all unique,
- the size of the matrices is not immediately related to the dimensionality of the system.

Therefore, the concept of Lax pairs does not provide a means to decide whether any given system is integrable (unless one is lucky to find a sufficiently large Lax pair).

Beisert: Lecture Notes on Integrability (p17)

## Applications:

- Classical mechanics (e.g. planetary motion)
- Classical field theories ( $1+1$ dimensions)
- Spin Chain Models
- $D=4 \mathrm{~N}=4 \mathrm{SYM}$ in the planar limit

Nonlinear Sciences > Exactly Solvable and Integrable Systems [Submitted on 12 Mar 2021]
Integrability ex machina

[^0]
## Formulating the search as optimisation

- Aim: Method to find new Lax pairs with unsupervised learning (i.e. not requiring prior knowledge of a Lax pair)
- Lax equation as loss:

$$
\dot{L}=[L, M] \rightarrow \mathscr{L}_{\mathrm{Lax}}=|\dot{L}-[L, M]|^{2}
$$

- Equivalence to EOM (e.g. $\dot{x}_{i}=f_{i}\left(x_{i}, \partial x_{i}, \ldots\right)$ ): $L$ has to include $x_{i}$ in some component (LHS of EOM), $[L, M]$ has to include RHS of EOM

$$
\begin{aligned}
\mathscr{L}_{\mathrm{L}} & =\sum_{i, j} \min _{k}\left(\left\|c_{i j k} \dot{L}-\dot{x}_{k}\right\|\left\|^{2},| | \dot{L}_{i j}\right\|^{2}\right)+\sum_{k} \min _{i j}\left(\left\|c_{i j k} \dot{L}_{i j}-\dot{x}_{k} \mid\right\|^{2}\right), \quad c_{i j k}=\frac{\sum_{\text {batch }} \dot{L}_{i j}}{\sum_{\text {batch }} \dot{x}_{k}} \\
\mathscr{L}_{\mathrm{LM}} & =\sum_{i, j} \min _{k}\left(\left\|\tilde{c}_{i j k}[L, M]_{i j}-f_{k}\right\|\left\|^{2},\right\|[L, M]_{i j} \|^{2}\right)+\sum_{k} \min _{i j}\left(\left\|\tilde{c}_{i j k}[L, M]_{i j}-f_{k} \mid\right\|^{2}\right), \tilde{c}_{i j k}=\frac{\sum_{b a t c h}[L, M]_{i j}}{\sum_{b a t c h} f_{k}}
\end{aligned}
$$

- Avoiding mode collapse:

$$
\mathscr{L}_{\mathrm{MC}}=\max \left(1-\sum\left|A_{i j}\right|, 0\right)
$$

- Total loss:

$$
\mathscr{L}_{\mathrm{Lax}-\mathrm{pair}}=\alpha_{1} \mathscr{L}_{\mathrm{Lax}}+\alpha_{2} \mathscr{L}_{\mathrm{L}}+\alpha_{3} \mathscr{L}_{\mathrm{LM}}+\alpha_{4} \mathscr{L}_{\mathrm{MC}}
$$

## Applications

## Harmonic Oscillator

- Harmonic Oscillator:

$$
H=\frac{1}{2} p^{2}+\frac{\omega^{2}}{2} q^{2} ; \quad \dot{q}=p, \quad \dot{p}=-\omega^{2} q
$$

- Lax Pair:

$$
L=\left(\begin{array}{cc}
0.437 q & -0.073 p \\
-0.666 p & -0.437 q
\end{array}\right), \quad M=\left(\begin{array}{cc}
0.001 & 0.329 \\
-3.043 & -0.001
\end{array}\right)
$$

- Consistency check:

$$
\frac{d L}{d t}=\left(\begin{array}{cc}
0.437 \dot{q} & -0.073 \dot{p} \\
-0.666 \dot{p} & -0.437 \dot{q}
\end{array}\right)=\left(\begin{array}{cc}
0.441 p & 0.288 q \\
2.660 q & -0.441 p
\end{array}\right)=[L, M]
$$

- Conserved quantities:

$$
L^{2}=\left(\begin{array}{cc}
0.048618 p^{2}+0.190969 q^{2} & 0 \\
0 & 0.048618 p^{2}+0.190969 q^{2}
\end{array}\right) \Rightarrow \operatorname{tr} L^{2} \approx 0.2 H
$$

## Applications

## Further systems

- Korteweg-de Vries (waves in shallow water):

$$
\dot{\phi}(x, t)+\phi^{\prime \prime \prime}(x, t)+6 \phi(x, t) \phi^{\prime}(x, t)=0
$$

- Heisenberg magnet:

$$
\begin{gathered}
H=\frac{1}{2} \int d x \vec{S}^{2}(x), \vec{S} \in S^{2} ; \text { constraint: } \\
\left\{S_{a}(x), S_{b}(y)\right\}=\epsilon_{a b c} S_{c}(x) \delta(x-y)
\end{gathered}
$$

- $\mathrm{O}(\mathrm{N})$ non-linear sigma models (Sine-Gordon equation and principal chiral model):

$$
\mathscr{L}=-\operatorname{Tr}\left(J_{\mu} J^{\mu}\right), \quad J_{\mu}=\left(\partial_{\mu} g\right) g^{-1}, \quad \mu=0,1 .
$$

## Perturbations on integrable systems

- Harmonic Oscillator:

$$
H_{0}=\frac{p_{x}^{2}+p_{y}^{2}}{2 m}+\omega^{2}\left(q_{x}^{2}+q_{y}^{2}\right)
$$

- Are the following perturbations integrable:

$$
H_{1}=\epsilon q_{x}^{2} q_{y}^{2}, \quad H_{2}=\epsilon q_{x} q_{y}
$$

- Initialise network at known solution for unperturbed system and see how it reacts to samples from perturbed
 system

Beyond symmetries, are there other structures in theoretical (particle) physics?

## Dualities

## Can they be useful in ML? <br> Can ML provide new perspectives on dualities?

## Dualities

## 2D Ising - Self-duality

Ordered rep. $\leftrightarrow$ Disordered rep.


## Field Theories

Electromagnetic Duality:

$$
\overrightarrow{\mathbf{E}} \leftrightarrow \overrightarrow{\mathbf{B}}
$$

el. charges $\leftrightarrow$ mag. monopoles

Seiberg Dualities in supersymmetric gauge theories:


## Holography



String Dualities
T-duality: winding \& momentum strings



## Essence of Dualities



## Essence of Dualities



## Essence of Dualities



## Connecting Dualities and Machine Learning

Does a neural network use such transformations automatically?


## Connecting Dualities and Machine Learning

Does a neural network use such transformations automatically?

If not, how can we make use of such transformations?

1) Latent loss to maximize distance between signal \& noise:

$\mathscr{L}=\max \left(0, \alpha-\xi_{1}^{2}-\xi_{2}^{2}\right)$,
$\xi_{i}^{2}$ largest square values of outputs




## Connecting Dualities and Machine Learning

Does a neural network use such transformations automatically?

If not, how can we make use of such transformations?
2) Pre-training with medium hard inference task on latent dimension

## Autoencoder



1D Ising Model with multiple spin-interactions: Feasible task: Energy
Hard inference task: Metastable state


Actual Dual variables


Intermediate variables

## Conclusions and Outlook

## Learning and using physics bias with ML

- Bias networks with physics knowledge for efficient results:
- Finding the functional bias possible: Learning mathematical structures (e.g. metric, Hamiltonian, symmetries) is possible in an unsupervised way when "appropriate" loss functions can be identified: - Symmetries from embedding layer without prior knowledge
- Symmetries from phase space samples
- Machinery for discovery of novel structures in integrability: Currently Lax pairs and connections for classical systems. Identify (some) integrable perturbations.
- Interpretation/enforcing of latent variables as variables of a dual theory (via appropriate losses)



## Thank you!

2012.04656: Numerical CY-Metrics
2104.14444: Simulations with Symmetry Control Neural Networks
2103.07475: Integrability
2003.13679: Symmetries from Embedding Layer

For talks at the interface of physics and ML: physicsmeetsml.org

## Control via Symmetries

- Losses to ensure appropriate functional forms:

$$
\begin{aligned}
& \mathscr{L}_{\mathrm{HNN}}=\sum_{i=1}^{N \cdot d}\left\|\frac{\partial \mathscr{H}_{\phi}(\mathbf{P}, \mathbf{Q})}{\partial p_{i}}-\frac{d q_{i}}{d t}\right\|_{2}+\left\|\frac{\partial \mathscr{H}_{\phi}(\mathbf{P}, \mathbf{Q})}{\partial q_{i}}+\frac{d p_{i}}{d t}\right\|_{2} \\
& \mathscr{L}_{\text {Poisson }}=\sum_{i, j=1}^{N \cdot d}\left\|\left\{Q_{i}, P_{j}\right\}-\delta_{i j}\right\|_{2}+\sum_{i, j>i}^{N \cdot d}\left\|\left\{P_{i}, P_{j}\right\}\right\|_{2}+\left\|\left\{Q_{i}, Q_{j}\right\}\right\|_{2} \\
& \mathscr{L}_{\mathrm{HQP}}^{(n)}=\sum_{i=1}^{n}\left\|\frac{d P_{i}}{d t}\right\|_{2}+\left\|\frac{d Q_{i}}{d t}-\frac{\partial \mathscr{H}_{\phi}(\mathbf{P}, \mathbf{Q})}{\partial P_{i}}\right\|_{2}+\beta \sum_{i=n+1}^{N \cdot d}\left\|\frac{d P_{i}}{d t}+\frac{\partial \mathscr{H}_{\phi}(\mathbf{P}, \mathbf{Q})}{\partial Q_{i}}\right\|_{2}+\left\|\frac{d Q_{i}}{d t}-\frac{\partial \mathscr{H}_{\phi}(\mathbf{P}, \mathbf{Q})}{\partial P_{i}}\right\|_{2}
\end{aligned}
$$


[^0]:    Sven Krippendorf, Dieter Lust, Marc Syvaeri

